

# Diffusion: Kolmogorov Equations, Feynman-Kac the Martingale Problem

Math653: SDEs and applications, Chapter 8

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# Overview

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## Review: Itô Diffusions & Generator

## Itô diffusion

$$dX_t = b(X_t) dt + \sigma(X_t) dB_t, \quad X_0 = x,$$

where  $B_t$  is  $m$ -dimensional, Lipschitz and linear growth.

## Generator: The Infinitesimal Expectation Operator

$$\mathbb{E}^x[f(X_h)] = f(x) + hAf(x) + o(h) \Leftrightarrow Af(x) = \lim_{h \rightarrow 0} \frac{\mathbb{E}^x[f(X_h)] - f(x)}{h}.$$

$$Af = b \cdot \nabla f + \frac{1}{2} a : D^2 f.$$

$$Af(x) = \sum_{i=1}^d b_i(x) \partial_i f(x) + \frac{1}{2} \sum_{i,j=1}^d a_{ij}(x) \partial_{ij} f(x).$$

**Derivation from Itô's Formula** For  $f \in C_b^2(\mathbb{R}^d)$ , Itô's formula gives

$$df(X_t) = \nabla f(X_t)^T b(X_t) dt + \frac{1}{2} a(X_t) : D^2 f(X_t) dt + \nabla f(X_t)^T \sigma(X_t) dB_t.$$

Integrating from 0 to  $h$  and taking expectation,

$$\mathbb{E}^x[f(X_h)] - f(x) = \mathbb{E}^x \left[ \int_0^h Af(X_s) ds \right].$$

Taking  $h \rightarrow 0$  gives the generator formula.

### Key point

The generator captures the local behavior of the diffusion by averaging over an infinitesimal time interval.

# Dynkin's Formula: The Accumulated Version

**Question:** How do infinitesimal changes accumulate over  $[0, \tau]$ ?

Let  $\tau$  be a bounded stopping time. Integrating Itô's formula gives

$$f(X_\tau) - f(x) = \int_0^\tau Af(X_s) ds + \int_0^\tau \nabla f(X_s)^T \sigma(X_s) dB_s.$$

## Dynkin's Formula

$$\mathbb{E}^x[f(X_\tau)] = f(x) + \mathbb{E}^x \left[ \int_0^\tau Af(X_s) ds \right].$$

- Generator: local (infinitesimal) behavior. Dynkin: global behavior.
- Class of  $f$ : from  $C_b^2(\mathbb{R}^d)$  to general  $C^2$ .
- Dynkin's formula bridges the SDE to PDEs and potential theory.

# Running Example: Brownian Motion with Drift

## Model

Let  $\mu \in \mathbb{R}$  and  $\sigma > 0$  be constants:

$$dX_t = \mu dt + \sigma dB_t, \quad X_0 = x,$$

Solution:  $X_t = x + \mu t + \sigma B_t$ .

Its generator is

$$Af(x) = \mu f'(x) + \frac{1}{2}\sigma^2 f''(x).$$

# Kolmogorov's Equations

# Backward Equation: What PDE Evolves Observables?

## Key question

How does  $u(t, x) = \mathbb{E}^x[f(X_t)]$  evolve in time?

Define the semigroup ( $P_{t+s} = P_t P_s$ ;  $P_0 = I$ )

$$P_t f(x) := \mathbb{E}^x[f(X_t)].$$

Then  $u(t, x) = P_t f(x)$  and  $u(0, x) = f(x)$ .

Formally,

$$\frac{u(t+h, x) - u(t, x)}{h} = \frac{P_t(P_h f - f)(x)}{h}.$$

By Dynkin's formula on the short interval  $[0, h]$ ,

$$P_h f(x) - f(x) = \mathbb{E}^x \left[ \int_0^h Af(X_s) ds \right] = hAf(x) + o(h).$$

$$\Rightarrow \boxed{\partial_t u(t, x) = Au(t, x), \quad u(0, x) = f(x).}$$

- The backward equation evolves test functions or observables
- $\partial_t P_t f = AP_t f$ ;  $P_t = e^{tA}$

**Example: Backward Equation in 1D**  $dX_t = b(X_t) dt + \sigma(X_t) dB_t$ , then

$$Af(x) = b(x)f'(x) + \frac{1}{2}\sigma^2(x)f''(x).$$

Hence

$$\partial_t u = b(x)\partial_x u + \frac{1}{2}\sigma^2(x)\partial_{xx} u, \quad u(0, x) = f(x).$$

For the running example  $dX_t = \mu dt + \sigma dB_t$ ,

$$\partial_t u = \mu\partial_x u + \frac{1}{2}\sigma^2\partial_{xx} u, \quad u(0, x) = f(x).$$

## Note

The drift contributes the first-order transport and the diffusion contributes the second-order term.

# Forward Equation: What PDE Evolves the Law?

## Key question

How does the probability density of  $X_t$  evolve?

Assume  $X_t^x$  has transition density  $p_t(x, y)$ :

$$P_t f(x) = \int_{\mathbb{R}^d} f(y) p_t(x, y) dy.$$

We want

- A PDE for  $p_t(x, y)$  in the variable  $y$ .
- A PDE for  $\rho(t, y)$ , the density of  $X_t$ , obtained by integrating over  $x$ .

Differentiate  $P_t f$  in time:

$$\int f(y) \partial_t p_t(x, y) dy = \partial_t P_t f(x) = P_t A f(x).$$

Rewrite the right-hand side by integration by parts in  $y$ :

$$P_t A f(x) = \int A f(y) p_t(x, y) dy = \int f(y) A_y^* p_t(x, y) dy.$$

[ $A^*$  is the adjoint operator; see next slide.]

The equation holds for all  $f \in C_c^2(\mathbb{R}^d)$ , so by calculus of variation,

$$\partial_t p_t(x, y) = A_y^* p_t(x, y).$$

- The forward equation is the dual version of the backward equation.
- It holds for each  $x$ . Integrating over  $x$ , we get the PDE for  $\rho(t, y)$ .

## Adjoint: The Dual Pairing Behind the Formula

In one dimension, let

$$Af = b(x)f'(x) + \frac{1}{2}a(x)f''(x), \quad a(x) = \sigma^2(x).$$

The dual action is defined by

$$\langle f, \rho \rangle := \int_{\mathbb{R}} f(x)\rho(x) dx.$$

Integration by parts gives

$$\langle Af, \rho \rangle = \left\langle f, -\partial_x(b\rho) + \frac{1}{2}\partial_{xx}(a\rho) \right\rangle.$$

Therefore,

$$A^*\rho = -\partial_x(b\rho) + \frac{1}{2}\partial_{xx}(a\rho).$$

# Backward vs. Forward: The Two PDE Viewpoints

## Backward PDE

Average observable  $f$ :

$$u(t, x) = \mathbb{E}^x[f(X_t)]$$

Equation:

$$\partial_t u = Au, \quad u(0, x) = f(x).$$

## Forward PDE

Density:

$$\rho(t, \cdot) = \text{law density of } X_t$$

Equation:

$$\partial_t \rho = A^* \rho, \quad \rho(0, \cdot) = \rho_0.$$

## Key point

The same diffusion gives two PDEs: one evolves observables, the other evolves probability mass.

**Running Example:**  $dX_t = \mu dt + \sigma dB_t$ . Its density evolves by

$$\partial_t \rho = A^* \rho = -\mu \partial_x \rho + \frac{1}{2} \sigma^2 \partial_{xx} \rho.$$

General 1D equation  $dX_t = b(X_t) dt + \sigma(X_t) dB_t$ , the density solves

$$\partial_t \rho = -\partial_x (b\rho) + \frac{1}{2} \partial_{xx} (a\rho).$$

- $-\partial_x (b\rho)$  transports mass by the drift,  $\frac{1}{2} \partial_{xx} (a\rho)$  diffuses mass.
- A heat equation with advection: the density is translated by the drift and spread by the Brownian noise.

### Further question

What can solutions of  $A^* \rho = 0$  tell us about invariant measures and long-time behavior?

# The Resolvent Operator

# Resolvent Operator: Compression in time

## Key question

The semigroup  $P_t$  tells us what happens at each time  $t$ . Can we compress all times into one operator?

The answer is yes: discount the future and integrate in time.

For  $\alpha > 0$ , define the resolvent

$$R_\alpha g(x) := \mathbb{E}^x \left[ \int_0^\infty e^{-\alpha t} g(X_t) dt \right] = \int_0^\infty e^{-\alpha t} P_t g(x) dt.$$

## Meaning

$R_\alpha g(x)$  is the expected total discounted payoff accumulated along the path. It is the Laplace transform of the semigroup.

**Resolvent Identity.** If  $f \in C_b^2$  and  $g = (\alpha - A)f$ , then Dynkin's formula applied to  $e^{-\alpha t}f(X_t)$  yields

$$\mathbb{E}^x \left[ \int_0^\infty e^{-\alpha t} (\alpha - A)f(X_t) dt \right] = f(x).$$

Thus

$$R_\alpha(\alpha - A)f = f.$$

Equivalently, if  $u = R_\alpha g$ , then formally

$$(\alpha - A)R_\alpha g = (\alpha - A)u = g.$$

- The resolvent  $R_\alpha$  is the Laplace transform of  $P_t = e^{tA}$  (formally):

$$R_\alpha = \int_0^\infty e^{-\alpha t} P_t dt = (\alpha - A)^{-1}.$$

- It solves an elliptic problem by damped-averaging along trajectories.

### Further question

How much spectral information about  $A$  is encoded in  $R_\alpha$ ?

# The Feynman-Kac Formula

# The Feynman-Kac Formula: solve PDE by SDE

## Key question

Can we solve a parabolic PDE by averaging over random paths?

Let  $q \geq 0$  and define

$$v(t, x) := \mathbb{E}^x \left[ \exp \left( - \int_0^t q(X_s) ds \right) f(X_t) \right].$$

Then  $v(0, x) = f(x)$  and, under standard regularity assumptions,

$$\partial_t v = Av - qv, \quad v(0, x) = f(x).$$

- When  $q = 0$ , we recover the backward equation.
- The potential term  $q$  discounts the future along the path.
- Example: If  $q \equiv \lambda > 0$ , then  $v(t, x) = \mathbb{E}^x [e^{-\lambda t} f(X_t)]$ .

## Theorem (Theorem 8.2.1)

Let  $f \in C_0^2(\mathbb{R}^n)$  and  $q \in C^1(\mathbb{R}^n)$  be bounded from below. Define

$$v(t, x) := \mathbb{E}^x \left[ \exp \left( - \int_0^t q(X_s) ds \right) f(X_t) \right].$$

Then:

- 1  $v$  solves the IVP

$$\partial_t v = Av - qv; \quad v(0, x) = f(x); \quad t > 0, x \in \mathbb{R}^n.$$

- 2 If  $w(t, x) \in C^{1,2}(\mathbb{R}_+ \times \mathbb{R}^n)$  is bounded on  $K \times \mathbb{R}_+$  for each compact  $K \subset \mathbb{R}^n$ , and if  $w$  solves the IVP, then  $w = v$ .

## Sketch of the proof.

**Part (1).** Let

$$Y_t := f(X_t), \quad Z_t := \exp\left(-\int_0^t q(X_s) ds\right).$$

Then  $dY_t = Af(X_t) dt + \nabla f(X_t)\sigma(X_t) dB_t$  and  $dZ_t = -Z_t q(X_t) dt$ . By Itô's product rule,

$$d(Y_t Z_t) = Z_t [Af(X_t) - q(X_t)f(X_t)] dt + Z_t \nabla f(X_t)\sigma(X_t) dB_t.$$

Then,  $v(t, x) = \mathbb{E}^x[Y_t Z_t]$  is differentiable in  $t$  and

$$\begin{aligned} \partial_t v(t, x) &= \lim_{h \rightarrow 0} \frac{\mathbb{E}^x[Y_{t+h} Z_{t+h}] - \mathbb{E}^x[Y_t Z_t]}{h} \\ &= \lim_{h \rightarrow 0} \frac{\mathbb{E}^x\left[\int_t^{t+h} Z_s [Af(X_s) - q(X_s)f(X_s)] ds\right]}{h} = Av - qv. \end{aligned}$$

Instead, use definition of the generator and the Markov property to get

$$\begin{aligned} Av(t, x) &= \lim_{r \rightarrow 0} \frac{1}{r} (\mathbb{E}^x[v(t, X_r)] - v(t, x)) \\ &= \lim_{r \rightarrow 0} \frac{1}{r} (\mathbb{E}^x [\mathbb{E}^{X_r}[Z_t f(X_t)]] - \mathbb{E}^x[Z_t f(X_t)]) \end{aligned}$$

Since  $e^{-\int_0^t q(X_{s+r}) ds} = e^{-\int_r^{t+r} q(X_s) ds} = Z_{t+r}/Z_r$  and the Markov property,

$$\begin{aligned} &\mathbb{E}^x [\mathbb{E}^{X_r}[Z_t f(X_t)]] - \mathbb{E}^x[Z_t f(X_t)] \\ &= \mathbb{E}^x [\mathbb{E}^x[f(X_{t+r})e^{-\int_0^t q(X_{s+r}) ds} \mid \mathcal{F}_r]] - \mathbb{E}^x[Z_t f(X_t)] \\ &= \mathbb{E}^x [\mathbb{E}^x[f(X_{t+r})Z_{t+r}/Z_r \mid \mathcal{F}_r]] - \mathbb{E}^x[Z_t f(X_t)] \\ &= \mathbb{E}^x [\mathbb{E}^x[f(X_{t+r})Z_{t+r} - f(X_t)Z_t \mid \mathcal{F}_r]] + \mathbb{E}^x[(Z_r^{-1} - 1)Z_{t+r}f(X_{t+r})] \\ &= \mathbb{E}^x [f(X_{t+r})Z_{t+r} - f(X_t)Z_t] + \mathbb{E}^x[(e^{\int_0^r q(X_s) ds} - 1)Z_{t+r}f(X_{t+r})] \\ &\frac{1}{r} \rightarrow \partial_t v(t, x) + q(x)v(t, x), \quad \text{as } r \rightarrow 0, \end{aligned}$$

where the second term converges to  $q(x)v(t, x)$  by the bounded convergence theorem.

**Proof of Part (2).** Consider extended process

$$H_t = (s - t, X_t^{0,x}, Z_t), \quad Z_t = z + \int_0^t q(X_r) dr,$$

and apply Dynkin's formula to  $\phi(H_t)$  with  $\phi(s, x, z) = e^{-z} w(s, x)$ .  
See the textbook for details.

## Generalization.

- **Feynman-Kac with time-dependent potential:**  $q = q(t, x)$ , and it can be even stochastic (fractional noise). Hu+Nualart+. 2008-2012.
- **Feynman-Kac with terminal-value.** For  $T > 0$  and  $g \in C_b^2(\mathbb{R}^d)$ ,

$$u(t, x) = \mathbb{E}^x \left[ e^{-\int_0^{T-t} q(X_s) ds} g(X_{T-t}) \right]$$

solves

$$\partial_t u + Au - qu = 0, \quad u(T, x) = g(x).$$

# The Martingale Problem

# Martingale Problem: weak solu. and mG characterization

## Key question

Characterize the law of a diffusion without fixing a Brownian motion?  
Weak solution to SDE v.s. generator with a martingale term?

For Ito diffusion  $dX_t = b(X_t)dt + \sigma(X_t)dB_t$  with  $X_0 = x$  (generator  $A$ )

$$M_t = f(X_t) - \int_0^t Af(X_s)ds = f(x) + \int_0^t \nabla f(X_s)^T \sigma(X_s)dB_s$$

is a martingale for any  $f \in C_0^2(\mathbb{R}^d)$ .

This only requires the existence of a weak solution, i.e., a probability measure  $\mathbb{P}^x$  on the path space such that the coordinate process  $\omega_t$  satisfies the SDE in distribution.

# Martingale Problem: Definition

## Definition

Let

$$Lf(x) = \sum_{i=1}^d b_i(x) \partial_i f(x) + \frac{1}{2} \sum_{i,j=1}^d a_{ij}(x) \partial_{ij} f(x),$$

where  $a$  and  $b$  are locally bounded Borel measurable;  $a(x) \geq 0$  (PSD) for all  $x$ . A probability measure  $\mathbb{P}^x$  on path space (all start from  $x$ ) solves the martingale problem for  $L$  if for every test function  $f \in C_c^\infty(\mathbb{R}^d)$ ,

$$M_t^f := f(\omega_t) - f(\omega_0) - \int_0^t Lf(\omega_s) ds$$

is a martingale in the canonical filtration.

The martingale problem is well-posed if there is a unique solution  $\mathbb{P}^x$ .

Strook and Varadhan (1979):

- weak solution  $\Leftrightarrow$  mG problem solution
- weak solution + Markov  $\Leftrightarrow$  well-posed mG problem

## Itô Processes vs. Itô Diffusions

# Itô Process and Itô Diffusion?

Diffusion:

$$\begin{aligned} \text{Itô diffusion:} & \quad dX_t = b(X_t) dt + \sigma(X_t) dB_t. \\ \text{Itô process:} & \quad dY_t = u(t, \omega) dt + v(t, \omega) dB_t, \end{aligned} \tag{1}$$

where  $u$  and  $v$  are adapted processes satisfying integrability conditions, but may depend on the whole past (not just the current state).

## Question

If  $X_t$  is an Itô process, then  $\phi(X_t)$  is also an Itô process for any  $\phi \in C^2$ .  
If  $X_t$  is an Itô diffusion, is  $\phi(X_t)$  also an Itô diffusion?

If its local characteristics depend only on the present state, then it may share the law of a diffusion.

### Theorem (Theorem 8.4.3: Itô Process vs. Diffusion)

Let  $X_t$  and  $Y_t$  be as in (1) with  $X_0 = Y_0 = x$ . Then,  $X_t$  has the same law as  $Y_t$  for each  $t \geq 0$  if and only if

$$\mathbb{E}[u(t, \omega) \mid \mathcal{N}_t^Y] = b(Y_t), \quad \mathbb{E}[v(t, \omega)v(t, \omega)^T \mid \mathcal{N}_t^Y] = \sigma\sigma^T(Y_t)$$

a.s. for each  $t \geq 0$ , where  $\mathcal{N}_t^Y$  is the  $\sigma$ -algebra generated by the path of  $Y$  up to time  $t$ .

The proof is based on the martingale problem and the following Brownian motion characterization.

### Theorem (Theorem 8.4.2)

An Ito process  $dY_t = vdB_t$ ;  $Y_0 = 0$  coincides in law with an  $n$ -dimensional Brownian motion if and only if  $vv^T(t, \omega) = I_n$  a.s. for each  $t \geq 0$ .

# Chapter 8 Takeaways

## Main messages

Path dynamics  $\rightarrow$  expectation  $\rightarrow$  generator  $\rightarrow$  backward/forward PDEs.  
Feynman-Kac: weighted paths average gives the solution to a PDE.

- The martingale problem characterizes the law.
- Ito process vs. diffusion.

## Big picture

Four equivalent lenses to study a diffusion: paths, generators, PDEs, and laws.